Mechanical characterization of sequentially layered photo-clickable thiol-ene hydrogels

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Article Info

Article history:
Received 2 July 2016
Received in revised form 2 September 2016
Accepted 4 September 2016
Available online 8 September 2016

Keywords:
Multi-layer hydrogel
Interfacial properties
Nanoindentation
Atomic force microscopy

Abstract

Multi-layer hydrogels are promising for tissue engineering due to the ability to control the local properties within each layer. However, the interface that forms between each layer has the potential to affect the performance of the hydrogel. The goals of this study were to characterize how the interface forms via its thickness and mechanical properties, identify its impact on the overall hydrogel properties, and provide new insights into how to control the interface. A photo-clickable poly(ethylene glycol) hydrogel was used to form bilayer hydrogels that were sequentially polymerized in a step-and-repeat process. Different processing conditions were studied: the time (0–20 min) before initiating polymerization of the second layer (soak time, \( t_s \)) and the hydrogel crosslink density (the same, less crosslinked, or more crosslinked) of the first layer as compared to the second layer. Interface thickness was characterized by confocal microscopy, monomer transport by Fickian diffusion, single and bilayer hydrogel mechanics by bulk moduli measurements, and interface moduli measurements using AFM, nanoindentation, and strain mapping. The interface thickness ranged from ~70 to 600 \( \mu \)m (1–10% of total height) depending on processing conditions, but did not affect the bulk hydrogel modulus. Analysis of monomer transport revealed that convection, due to changes in hydrogel swelling, and diffusion contribute to interface thickness. Nanomechanical analysis of bilayer hydrogels formed from soft (75 kPa) and stiff (250 kPa) layers showed a gradient in elastic modulus across the interface, which corresponded to strain maps. In summary, this work identifies that diffusive and convective transport of monomers across the interface controls its thickness and that a mechanically robust interface forms, which does not affect the hydrogel modulus. By controlling the processing conditions, the thickness of the interface can be tuned without affecting the mechanical properties of the bulk hydrogel.

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1. Introduction

Many tissues in the body exist as multi-layer structures with each layer having distinct biochemical and mechanical properties that together contribute to the overall function of the tissue. Examples of such tissues include cartilage, which is composed of superficial, middle and deep zonal layers, skin, which consists of an epidermis and dermis layer, and the junction that connects cartilage to bone or tendons and ligaments to bone. Between each layered structure, an interface forms that can be characterized by either extended or abrupt gradients in the biochemical and/or mechanical properties. For example, abrupt gradients are present in the mechanical transition at the osteochondral junction between stiff calcified cartilage and compliant hyaline cartilage (Campbell et al., 2012) and in the compositional transition across the muscle-tendon junction (Ker, 2007). The ability to mimic these complex tissue structures in a scaffold is an important step towards developing successful strategies for tissue engineering composite tissues.

Multi-layered hydrogels can be fabricated such that their chemistry and properties (e.g., crosslinking and porosity) vary in space. Mild polymerization conditions coupled with a high water content make hydrogels particularly suitable for the encapsulation and culture of cells in 3D scaffolds (Nicodemus and Bryant, 2008).
As such, multi-layer hydrogels are being designed and investigated for tissue engineering of complex multi-structured tissues. For example, multi-layer fibrous hydrogels made from chitosan were designed where keratinocytes were seeded into a layer simulating the dermis, fibroblasts seeded into a second layer representing the protective epidermis, and a thin film to mimic the basement membrane to engineer a skin substitute (Lin et al., 2015). This tri-layered hydrogel enabled both cell types to replicate the striation of full thickness skin more accurately and in a shorter time than single or bilayer counterparts. In a separate study, a multi-layer hydrogel formed from crosslinked poly(ethylene glycol) (PEG) was tuned by introducing different extracellular matrix molecules into each layer to mimic the chemical makeup of the different zones in cartilage (Nguyen et al., 2011). In doing so, 3D hydrogel niches were created that allowed a single population of stem cells to differentiate into zone-specific chondrocytes. In another study, the composition and mechanical properties in each layer of a cross-linked PEG hydrogel was varied to capture the physiochemical cues that arise in osteochondral tissues under mechanical loads (Steinmetz et al., 2015). When human mesenchymal stem cells were encapsulated in the multi-layer hydrogel and subjected to dynamic loading, differentiation towards a chondrogenic phenotype was observed in one layer and an osteogenic phenotype was observed in the other layer. Taken together, these examples among others demonstrate that multi-layered hydrogels are promising for tissue engineering complex multi-layer composite tissues.

In forming multi-layer hydrogels, the interface that forms between two adjacent layers is critical to the overall function of the hydrogel. For example, if the layers are not well integrated and secured together, the interface is prone to shear failure and can lead to delamination (Kandel et al., 2006; Sherwood et al., 2002). For interfaces with abrupt gradients, the two dissimilar materials disrupt strain transfer and can lead to areas of high stress concentration that result in failure at the interface (Erdogan, 1995). On the contrary, if the interface that forms is large, the local biochemical and mechanical properties can influence the fate of cells embedded within this region (Steinmetz et al., 2015). Thus, controlling the interface is important to designing multi-layer hydrogels that recapitulate extended or abrupt gradients found in native interfaces and that lead to mechanically robust interfaces to minimize failure. The resulting interface when multi-layer hydrogels are formed has received little attention.

This study investigates a bilayer PEG hydrogel and the interface that forms between the two layers. PEG hydrogels were chosen for their promise in producing multi-layer hydrogels for use in vascular (Fischer et al., 2015; Shinohara et al., 2013) and musculoskeletal (Fuoco et al., 2012; Hwang et al., 2010; Lin et al., 2014; Nguyen et al., 2012; Paxton et al., 2009; Steinmetz et al., 2015) tissue engineering. The objectives for this study are to characterize how the interface forms via its thickness and mechanical properties, identify its impact on the overall hydrogel properties, and provide new insights into how to control the interface. Bilayer PEG hydrogels were sequentially polymerized from photo-clickable thiol-ene macromolecular monomers. This reaction scheme was chosen for its promise in tissue engineering (Fairbanks et al., 2009; Tibbitt and Anseth, 2009) and its highly specific, efficient, and rapid reaction (Hoyle and Bowman, 2010). Specifically, this study investigated (a) the effects of monomer transport between the two layers as a function of processing conditions on interface thickness and bulk mechanical properties and (b) the local mechanical properties and strain transfer across the interface when each hydrogel layer is formed with different mechanical properties. The latter is important for applications where mechanical forces are applied to multi-layer hydrogels. Overall, this study provides new insight into how processing parameters influence the formation and properties of the interface and ultimately the contribution that the interface has on the macroscopic properties. Findings from this study will aid the development of multi-layer hydrogels where the interface can be tuned to create thin interfaces that lead to abrupt property changes across the interface or thick interfaces where large property gradients across the interface are required.

2. Materials and methods

2.1. Monomer synthesis

8-arm PEG with terminal amines (20,000 g/mol; JenKem Technology USA, Plano, TX) was functionalized with norbornenes by reacting 5-norbornene-2-carboxylic acid (Sigma-Aldrich, St. Louis, MO) with 2-(1H-7-azabenzo triazol-1-yl)-1,3,3-tetramethyl uran hexafluorophosphate methanaminium (HATU) (Chem-Impex International, Inc., Wool Dale, IL), and N,N-dioisopropyl-ethylthiylamine (DIPEA) (Chem-Impex) in dimethylformamide (DMF)/ dichloromethane (DCM) (Sigma-Aldrich). The reaction was allowed to proceed overnight at room temperature under an inert atmosphere. The product was precipitated in diethyl ether (Sigma-Aldrich), filtered, dialyzed, and lyophilized. The extent of conjugation of norbornene to each arm of the 8-arm PEG-amine was determined to be 92% using 1H NMR by comparing the protons across the carbon-carbon double bond in the norbornene to the methylene protons in PEG.

2.2. Fabrication of PEG hydrogels

A precursor solution consisting of 8-arm PEG-norbornene (PEGnor) monomer, PEG-dithiol (PEGdt) crosslinker (1000 or 3400 g/mol; Sigma-Aldrich) at 1:1 thiol:ene ratio (assuming 100% norbornene conjugation), and 0.05% (w/w) photoinitiator, 1-(4-(2-Hydroxyethyl)-phenyl)-2-hydroxy-2-methyl-1-propane-1-one (12959; Ciba Specialty Chemicals, Tarrytown, NY), in deionized water (diH2O) was photopolymerized by ultraviolet light (10 min, 10 mW/cm2, 352 nm). The concentration of 8-arm PEG-norbornene varied by 10, 15 or 25% (w/w) (Hydrogel I, II/III and IV, respectively) depending on the experiment. Hydrogels were formed in either a cylindrical or rectangular geometry depending on the experiment (as described below). For single layer constructs, full height (i.e., 5 mm) hydrogels were fabricated from a single macromer solution. For bilayer constructs, half-height (i.e., 2.5 mm) hydrogels were fabricated from one macromer solution that was polymerized to form the first layer. Previous studies have reported that complete polymerization of similar precursor formulations is reached in less than a minute (Roberts and Bryant, 2013) and therefore ten minutes ensures that the first layer has polymerized completely. Following polymerization of the first layer, a second macromer solution was carefully deposited on top of the first layer to form the second layer and to reach full height (i.e., 5 mm). The second solution was left to allow for transport of the monomers into the first hydrogel layer for prescribed periods of time (t), after which the second layer was polymerized. A schematic representation of the fabrication process is shown in Fig. 1A.

2.3. Swelling studies and hydrogel characterization

Single layer, full height cylindrical hydrogels prepared as stated above were weighed immediately after polymerization and then allowed to swell in diH2O for 48 h at room temperature. Equilibrium swollen hydrogels (n = 3/group) were weighed to determine the equilibrium swollen mass. The dry polymer mass was obtained for each hydrogel after lyophilization. The equilibrium mass swelling ratio, q, was calculated by dividing the equilibrium swollen mass by the polymer dry mass. The equilibrium swelling
ratio, $Q$, was estimated using densities of the respective polymer and solvent by

$$Q = 1 + \frac{\rho_p}{\rho_s} (q - 1).$$

Here $\rho_p$ is the density of the polymer, which for PEG is assumed 1.07 g/mL, and $\rho_s$ is the density of the solvent, which in this case for water is 1 g/mL. The hydrogel mesh size, $\xi$, was calculated by the following equation (Canal and Peppas, 1989),

$$\xi = \frac{v_{\lambda,2}^{1/3} C_p^{1/3} n^{1/2}}{Q^{1/2}},$$

where $v_{\lambda,2}$ is the equilibrium polymer volume fraction in the gel (i.e., inverse of the swelling ratio, $Q$), $\ell$ is the average bond length, which for PEG is $l = 1.49 \lambda$, $C_p$ is the characteristic ratio of the polymer, which is assumed to be 4.0 (Merrill et al., 1993), and $n$ is the number of bonds between the crosslinks, which was determined from the crosslink density. The crosslink density and the polymer-solvent interaction ($\chi_{12}$) parameter were estimated using a self-learning algorithm (Akalp et al., 2015) that combines Flory–Rehner theory with theories of mixture and poroelasticity (Flory and Erman, 1982; Flory, 1953; Holmes and Mow, 1990; Peppas, 1986). The crosslink density was determined to be 0.020, 0.026, 0.032, and 0.059 M and $\chi_{12}$ was determined to be 0.488, 0.451, 0.489, and 0.486 for Hydrogels I, II, III, and IV respectively.

### 2.4. Interfacial thickness measurement

Rectangular PEG hydrogels (5 mm height, 5 mm width, 1 mm thick) were made following methods described above. Fluorescently labeled hydrogels were formed by incorporating a fluorophore into each layer. In brief, 0.01 mM Alexa Fluor® 546 C5 Maleimide (Thermo Fisher Scientific, Waltham, MA) was added to one precursor solution and 0.01 mM Alexa Fluor® 488 C5 Maleimide (Thermo Fisher Scientific) was added to the second precursor solution prior to polymerization leading to the formation of a dual-fluorescent bilayer hydrogel. Hydrogels were imaged by confocal microscopy (Zeiss LSM 5 Pascal system using a Zeiss Axiovert microscope). Interface thickness was determined from confocal microscopy images and NIH Image J software. In brief, a line was manually drawn perpendicular to the interface and intensity along the distance of the line was determined for each of the fluorophore channels. The total distance across which both fluorophore channels detected fluorescence signal over the background denoted the interface thickness.

### 2.5. Mechanical testing of PEG hydrogels

Cylindrical PEG hydrogels (5 mm height and 5 mm diameter) were made following methods described above. Compressive testing was conducted on a Mechanical Testing System (MTS; Eden...
Prairie, MN). Equilibrium swollen hydrogels (n=3/group) were tested in unconflated compression on dry plates with minimal sticking at a constant strain rate of 0.5 mm/min to 50% strain (2 mN pre-load; limited by a 9 N max load with a total load cell capacity of 10 N). The tangent compressive modulus was calculated from the slope of the stress–strain curve at low (10–15%) strain and high (40–45%) strain.

2.6. One-dimensional diffusion studies

Rectangular PEG hydrogels (10 mm width, 25 mm length, 1 mm thick) were made between two glass slides following polymerization methods described above. The hydrogel was maintained between the two glass slides and reservoirs on either end of the hydrogel were filled with diH₂O. Hydrogels were allowed to swell for 48 h at room temperature. Pre-swollen hydrogels were placed in a VersaDoc MP 4000 Molecular Imaging System (Bio-Rad; Hercules, CA) and one reservoir was replaced with either an aqueous solution of 0.01 mM Alexa Fluor® 488 C5 Maleimide pre-reacted with 10 mM 8-arm PEG-thiol (20,000 g/mol; JenKem Technology USA, Plano, TX) or 0.01 mM Alexa Fluor® 488 C5 Maleimide pre-reacted with 100 mM methoxy PEG-thiol (800 MW; Sigma-Aldrich). Images were acquired at 12, 39, and 46 h for the 8-arm PEG molecule and at 24 and 30 h for the methoxy PEG-thiol molecule. From these images, Image J was used to plot concentration profiles of fluorescence intensity as a function of distance into the gel. These profiles were fit to the solution of Fick’s second law in one dimension via the curve fitting tool in MATLAB (The Mathworks, Inc., Natick, MA) to extract approximate diffusivities for the hydrogel conditions explored.

2.7. Atomic force microscopy

Rectangular PEG hydrogels (5 mm height, 5 mm width, 1 mm thick) were made following methods described above. An atomic force microscope (MFP-3D Classic, Santa Barbara, CA) was used to measure the spatially varying elastic moduli of both the single layer and bilayer hydrogels. Briefly, force spectroscopy (force-volume mapping) was performed in deionized water at 25 °C in a fluid cell for the MFP-3D AFM. For the single layer hydrogels, the measurements were carried out at each node of a 64 × 64 grid over a 10 μm × 10 μm area, while for the bilayer hydrogels, the measurements were carried out at each node of a 512 × 64 grid over an 80 μm × 10 μm area. Bruker SNL-10-B AFM probes (Cammarillo, CA) with triangular SiN cantilevers and Si tips were used. Each probe was calibrated with the thermal fluctuation method (Hutter and Bechhoefer, 1993); resulting values for the spring constant ks varied from 0.08 N/m to 0.16 N/m. Moreover, each probe was inspected before and after force spectroscopy with scanning electron microscopy (SEM) to assess the half-angle α and check for damage or material transfer; α varied from 15° to 23° and minimal material transfer was observed. AFM force-displacement (F–d) data was acquired with a peak load of 20 nN at a displacement rate of 10 μm/s. Force-deformation (F–δ) data were derived from the raw F–d curves by subtracting the cantilever deflection F/kc. To assess Young’s modulus E at each node, the loading portion of each F–δ curve was fit to an analytical model for a rigid conical tip in contact with an elastic half-space (Sneddon, 1965), i.e., \( F = (2/π)(E/(1–ν^2))(tan α)δ^2 \) using an average value for α from SEM, an assumed Poisson’s ratio for the hydrogel (ν = 0.5), and E as the sole fitting parameter.

2.8. Nanoindentation

Rectangular PEG hydrogels (5 mm height, 5 mm width, 1 mm thick) were made as described above for nanoindentation testing. A Ti-950 Triboindenter (Hysitron, Minneapolis, MN) performed micro scale indentation of single layer and bilayer hydrogels. Five samples were submerged in deionized water at 27 °C and a fluid cell 100 μm radius cono-spherical probe was pressed into the samples to a depth of 10 μm at a rate of 2 μm/s. Each sample was indented 300 times in two interlocking arrays. The first array was 20 indents wide by 8 indents tall with 100 μm spacing in x and 200 μm spacing in y. The second array had the same spacing, but was 20 indents wide by 7 indents tall with an offset from the starting location of the first array by 50 μm in x and 100 μm in y. To avoid effects from adjacent tests, all indents were spaced at three times the contact radius. Because the material behaved elastically the reduced modulus at each indent location was found using the maximum force and depth for each load displacement curve. The testing was performed on an extended displacement stage (Hysitron, xZ 500 Extended Displacement Stage), the load function was trapezoidal (5 s load, 30 s hold and 5 s unload), the Poisson ratio was assumed to be 0.5 for the samples, and the Hertz analytical method was used to evaluate the elastic modulus (Fischer-Cripps, 2002; Hertz, 1881).

2.9. Bead tracking and image processing

Cylindrical PEG hydrogels (5 mm height and 5 mm diameter) were made following methods described above. Silicon dioxide microparticles (5 μm; Sigma-Aldrich) were entrapped into each hydrogel layer at a concentration of 0.3% solids. The density of microparticles was chosen to be low enough so as not to alter the mechanical properties of the hydrogel, while enabling visualization of multiple beads in the hydrogel by a bright field microscope (Zeiss Axiovert 40C). To minimize microparticle settling, the macromer solution was vortexed immediately before polymerization. To distinguish between the layers, PolyFlour 570 (Methacryloyloxyethyl Thiocarbonyl Rhodamine B; Polysciences Inc., Warrington, PA) was added to the precursor solution for the second layer creating a red tint. These samples were then cut in half with a razor blade, placed cut side down on a windowed compression rig, loaded in compression, and images taken at 0.5% strain increments. A MATLAB based Digital Image Correlation routine (Jones, 2013; Jones et al., 2014) was used to track the movement of the microparticles in the direction of the applied strain and extract hydrogel deformation. These deformations were then applied to a two dimensional (2D) finite element model to compute sample strain maps for each image.

2.10. Statistics

Statistics were performed using KaleidaGraph 4.1.3 software (Synergy Software, Reading, PA). Significant differences were established using a one-way Analysis of Variance with soak time as a factor. Comparisons between bilayer hydrogels and single layer hydrogels were performed by an unpaired, equal variance two-sample t-test. P-values are provided to indicate significance. All numerical results are presented as mean (standard deviation) and graphical results are presented as mean with standard deviation as error bars (n=3).

3. Results

3.1. Fabrication of bilayer hydrogels

Bilayer hydrogels were formed with covalently tethered fluorophores to enable visualization of each layer and the interface. The precursor solution for the first layer contained a red fluorophore and was polymerized. The precursor solution for the second layer
contained a green fluorophore and was gently applied on top of the first layer. The time prior to exposing the second layer to light, herein referred to as soak time \( t_s \), was varied from 0 to 20 min. A schematic of the process is shown in Fig. 1A. The process resulted in interfaces that were visualized by the overlap of the two fluorophore channels (Fig. 1B). Four different hydrogel formulations, which are referred to as Hydrogels I–IV, with varying properties from soft to stiff were investigated to form the bilayer hydrogels. The formulation and resulting properties for each hydrogel formulation are given in Fig. 1C. The equilibrium swelling of hydrogels. The formulation and resulting properties for each hydrogel formulation are given in Fig. 1C. The equilibrium swelling ratio varied from 18 for Hydrogel II to 10 for Hydrogel IV, while the mesh size varied from 42 to 23 nm, respectively. The compressive modulus also varied from 75 to 250 kPa for hydrogel conditions (Fig. 1C). Three bilayer hydrogels were fabricated from the different hydrogel formulations to form Bilayer III:III, Bilayer II:III, and Bilayer IV:1, where the first Roman numeral corresponds to the hydrogel formulation in the first layer and the second Roman numeral corresponds to the hydrogel formulation in the second layer (Fig. 1D). Representative photographs for each bilayer condition after reaching equilibrium swelling are shown in Fig. 1E.

3.2. Increased time for transport leads to larger interfaces but no effect on bulk compressive modulus

Bilayer III:III was used to investigate the effect of soak time \( t_s \) on the resulting interface and macroscopic properties in the form of compressive modulus. The interface thickness increased \((p<0.0001)\) with soak time, for example from 65(14) μm with a \( t_s \) of 0 min to 230(17) μm with a \( t_s \) of 20 min (Fig. 2A–C). The interface thickness was \(~1–4\%\) of the total hydrogel height (Fig. 2C).

3.3. An increase in mesh size leads to larger interfaces and affects the bulk compressive modulus

Bilayer II:III was used to investigate the effect of a larger mesh size in the first layer on the resulting interface thickness and compressive modulus. The interface thickness increased \((p<0.0001)\) with increasing soak time, for example from 130(22) μm with \( t_s = 0 \) min to 580(5) μm with \( t_s = 20 \) min (Fig. 3A). The overall thickness of the interface was substantially larger in this bilayer hydrogel compared to Bilayer III:III. For example, at \( t_s = 20 \) min the interface thickness was 2.5-fold larger in Bilayer II:III compared to Bilayer III:III. The interfacial thickness was \(~2–10\%\) of the total hydrogel height (Fig. 3C). The compressive modulus of Bilayer II:III was not significantly affected by soak time at the low \((p=0.13)\) or high \((p=0.21)\) strains with no observable trends (Fig. 3C). At low strains, the compressive modulus of the bilayer hydrogels for all soak times was either lower or not different \((p=0.006–0.49)\) when compared to the single layer Hydrogel II, but was lower \((p=0.001–0.07)\) when compared to the single layer Hydrogel III (i.e., the more highly crosslinked hydrogel) (Fig. 3D). At high strains, the compressive modulus of the bilayer hydrogels for all soak times was higher \((p=0.038–0.14)\) than Hydrogel II, and

Fig. 2. (A) Representative confocal microscopy images of Bilayer III:III, fabricated with a bottom layer from 15% (w/w) PEGnor and PEG4kdt (red), a top layer from 15% (w/w) PEGnor and PEG1kdt (green), and the interface (yellow); scale bar =200 μm. (B) A representative plot of intensity for green and red fluorescence across the interface for Bilayer III:III (C) Interface thickness and average percent height of the interface relative to the total height of the hydrogel as a function of soak time \( t_s \) for Bilayer III:III. (D) Compressive modulus at low (10–15%) and high (40–45%) strains of the Bilayer III:III hydrogel and single layer Hydrogel III. Data are presented as mean \((n=3)\) with standard deviation as error bars; \( p \)-values in panel D denote significance from single layer hydrogel. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)
strains and affected (from 140 the compressive modulus. The mean interface thickness ranged mesh size in the bulk compressive modulus. Smaller mesh size does not affect interface thickness, but affects signs of failure up to 45% strain.

3.4. Smaller mesh size does not affect interface thickness, but affects the bulk compressive modulus

Bilayer IV:1 was used to investigate the effects of a smaller mesh size in the first layer on the resulting interface thickness and the compressive modulus. The mean interface thickness ranged from 140–210 μm with soak time, but there were no observable trends (p=0.013) (Fig. 4A). The interface thickness was ~3–4% of the total hydrogel height (Fig. 4C). The compressive modulus of the bilayer hydrogels was affected (p=0.0005) by soak time at low strains and affected (p=0.006) by soak time at the high strains, but there were no observable trends. At low strains, the compressive modulus of the bilayer hydrogels for all soak times was higher (p=0.0002–0.27) than single layer Hydrogel I, but lower (p=0.0003–0.0008) than single layer Hydrogel IV (Fig. 4D). At high strains, the compressive modulus of the bilayer hydrogels for all soak times was higher (p=0.0001–0.0005) than Hydrogel I, but lower (p=0.0005–0.0037) than Hydrogel IV. None of the bilayer conditions delaminated or showed any observable signs of failure up to 45% strain.

3.5. Smaller mesh size inhibits diffusion of monomer into the first hydrogel layer

When the second precursor solution is applied to the already polymerized first layer, the monomers from the precursor solution will diffuse into the first layer hydrogel. To characterize this diffusion, diffusive transport of a fluorescently labeled version of each monomer, the 8-arm PEG and the PEG crosslinker, into Hydrogel II, III, and IV was measured. A one-dimensional experimental set-up shown in Fig. 5A was used to experimentally measure concentration as a function of distance and time. The molecular weight of each fluorescently labeled molecule was similar to the corresponding monomer. To determine the diffusion coefficient for each monomer in each of the hydrogels, Fickian diffusion was assumed and the experimental data were fit to

\[ C(x, t) = C_0 \text{erf} \left( \frac{x}{2\sqrt{Dt}} \right) \]

where \( C \) is concentration of the fluorescently labeled PEG molecule, \( C_0 \) is the concentration in the reservoir, which was held constant, \( D \) is the diffusion coefficient, \( x \) is distance and \( t \) is time. This equation was determined by solving the simple case of diffusion with time \( t \) in one dimension (taken as the \( x \)-axis) with initial and boundary conditions \( C(x,0)=0 \) and \( C(0,t)=C_0 \). Select data are shown in Fig. 5B along with the estimated diffusion coefficients for the fluorescently labeled PEG monomers in each of the three hydrogels (Fig. 5C). As expected, hydrogels with a smaller mesh size resulted in lower diffusivities for both monomers. In addition, due to its size, the smaller monomer, the PEG crosslinker, diffused to a greater extent as compared to the larger 8-arm PEG monomer. Using these diffusivities, the concentration profile of each monomer as it diffuses into the first layer, can be estimated for each bilayer condition. The initial concentration was assumed to be the concentration of the respective monomer in the precursor solution for the second hydrogel layer, where Hydrogel III had a 34 mM PEGdt and 8.1 mM PEGnor concentration, and remained lower (p=0.002–0.01) than Hydrogel III (Fig. 3D). None of the bilayer conditions delaminated or showed any observable signs of failure up to 45% strain.
Hydrogel I had a 21 mM PEGdt and 5.1 mM PEGnor concentration. The concentration profile of each monomer into the first layer was plotted at a time of 20 min, corresponding to the \( t_s = 20 \) min condition, and is shown in Fig. 5D. The interface that was experimentally measured in the bilayer hydrogels for the \( t_s = 20 \) min condition (i.e., in Figs. 2C, 3C, and 4C) was superimposed on each plot. This overlay can then be used to predict, based on Fickian diffusion, the concentration of the PEG monomers within the interfacial layer. For Bilayer II:III, the concentration of PEGdt ranged from 34 mM to 2.9 mM and that of the PEGnor ranged from 8.1 mM to 0.006 mM. Finally, for Bilayer IV:1, the concentration of PEGdt ranged from 21 mM to 2.4 mM and that of the PEGnor ranged from 5.1 mM to 0.01 mM.

3.6. Spatial mapping of local mechanical properties reveals a gradual transition in elastic modulus at the interface of a bilayer hydrogel

To characterize the local mechanical properties across the interface, nanoindentation and AFM were performed. When used together, nanoindentation is capable of mapping modulus over large, millimeter-sized 2D surfaces and AFM facilitates measurement of modulus variations across small, sub-millimeter regions with nanometer-scale spatial resolution. To evaluate how hydrogel composition influenced the properties of the interface, we selected Bilayer IV:1 because it exhibited the largest difference in compressive modulus between the two layers. Spherical silicon dioxide beads (5 \( \mu \)m in diameter), encapsulated within hydrogels, are much larger than the mesh size of the hydrogels and therefore their movement can be correlated to the local strain of the hydrogel. Digital image correlation methods and a 2D model were used to calculate finite strains across the bilayer hydrogel. Under a 15% strain applied to the bilayer hydrogel, a color map of the strain distribution within the hydrogel was generated (Fig. 7A). The soft layer experienced \( \approx 20\% \) strain while counterparts of each hydrogel, Hydrogel I and IV, was first evaluated (Fig. 6A) and the resulting values largely agree with bulk compressive modulus measurements at 10–15% strains (Fig. 4D). In the bilayer hydrogel, AFM demonstrated a gradient in the modulus from the stiff Hydrogel IV layer to the soft Hydrogel I layer that spanned \( \approx 70 \mu \)m of the 80 \( \mu \)m-long scan region (Fig. 6B). Nanoindentation was also used to evaluate moduli of the individual layers and the interface region. Nanoindentation moduli, collected at low (\( \approx 4\% \)) strains when tested using a large radius conospherical probe produced correspondingly lower modulus values for the hydrogels, as compared to values from AFM and unconfined compression (Fig. 6C). Nanoindentation also demonstrated a graded transition of indentation modulus from the stiff to the soft layer (Fig. 6C).

3.7. Strain transfer depends on the local mechanical properties of each layer

Finite strains within Bilayer IV:1 were determined by bead tracking and digital image correlation techniques under compressive loading. Bilayer IV:1 was chosen because it exhibited the largest difference in compressive modulus between the two layers. Spherical silicon dioxide beads (5 \( \mu \)m in diameter), encapsulated within hydrogels, are much larger than the mesh size of the hydrogels and therefore their movement can be correlated to the local strain of the hydrogel. Digital image correlation methods and a 2D model were used to calculate finite strains across the bilayer hydrogel. Under a 15% strain applied to the bilayer hydrogel, a color map of the strain distribution within the hydrogel was generated (Fig. 7A). The soft layer experienced \( \approx 20\% \) strain while...
the stiff layer experienced \( \sim 10\% \) strain. The strains were plotted as a function of position for increasing bulk strains (Fig. 7B). Across the interface, the strain levels progressively decreased from the soft to the stiff layer. The variation in the strain from the soft to stiff layers occurred over a relatively large length scale of \( \sim 1 \) mm.

4. Discussion

The overall objectives for this study were to characterize how the interface forms via its thickness and mechanical properties, identify its impact on the overall hydrogel properties, and provide new insights into how to control the interface when multi-layer hydrogels are formed from sequential polymerization of monomers in a step-and-repeat process. Several factors controlled the thickness of the interface and included soak time (i.e., time before the second layer is polymerized) and the mesh size of the first layer. The interface thickness ranged from \( \sim 70 \) to 600 \( \mu \)m depending on these factors, but did not influence the bulk modulus measurements. We further showed that the interface results in a gradient in mechanical properties when bilayer hydrogels are formed from soft and stiff layers and this translated to a gradient in strain transfer across the interface under a gross applied strain. Overall, the study shows that a mechanically robust interface forms, but whose thickness can be controlled through the choice of processing conditions.

In the step-and-repeat process, the second hydrogel precursor solution is deposited on top of a crosslinked hydrogel immediately after it is polymerized. There are two potential factors that can influence the formation of the interface and its thickness: (1) transport by diffusion of monomers into the first hydrogel layer (Amsden, 1998a, 1998b; Hagel et al., 2013; Lustig and Peppas, 1988) and (2) transport by convection of monomers due to changes in swelling of the first hydrogel layer (Brannon-Peppas and Peppas, 1991; Kim et al., 1992; Martens and Anseth, 2000). Monomer transport by diffusion will occur due to a concentration gradient, but will depend on the relative size of the monomers in the second layer to the mesh size of the hydrogel in the first layer.

**Fig. 5.** (A) Schematic of the experimental setup for 1D diffusion for each monomer (i.e., a methoxy-PEG-fluorophore representing the PEG\(_{\text{dt}}\) crosslinker or a fluorescently labeled 8-arm PEG representing the PEG\(_{\text{nor}}\) monomer) into each of the hydrogels used in the first layer of the bilayer hydrogels. Hydrogels were sandwiched between two glass slides and a silicone mold (gray) with reservoirs on each end. Green represents the fluorescently labeled 8-arm PEG or fluorescently labeled PEG crosslinker in solution in the reservoir and diffusing into the hydrogel. (B) Representative plots for the 1D concentration profile of fluorescently labeled PEG monomers as a function of distance into a hydrogel. Data for Hydrogel II are shown with the fluorescently labeled PEG crosslinker representing PEG\(_{\text{dt}}\) at the 30 h timepoint and with the fluorescently labeled 8-arm PEG representing the PEG\(_{\text{nor}}\) monomer at the 39 h timepoint. Dotted line represents experimentally determined concentration profiles. Solid line represents the accompanying fit to the solution of Fick’s second law in 1D, which was used to determine the diffusion coefficient for each hydrogel formulation and shown in (C) with mean and (standard deviation). (D) The predicted concentration profiles at \( t = 20 \) min for each hydrogel. Dashed line represents the concentration of PEG\(_{\text{dt}}\) as a function of distance and the dashed and dotted line represents that of the PEG\(_{\text{nor}}\) monomer. Superimposed in gray is the experimentally determined interface as measured by confocal microscopy with a dot representing the intersection of the edge of the interface and the concentration of PEG\(_{\text{dt}}\) due to transport by diffusion, and a square representing the intersection of the edge of the interface and the concentration of PEG\(_{\text{nor}}\) due to transport by diffusion.
when the first layer was fabricated from Hydrogel II. The relatively large interface of \( \sim 130 \mu m \) at the initial time point suggests transport by convection due to a rapid initial swelling response. Indeed, this hydrogel swelled 3.9x its initial mass when placed in an aqueous solvent, which can be seen in Fig. 1E. With longer soak times, the interface continued to grow. The relatively large mesh size of Hydrogel II led to the highest diffusivity for the PEG crosslinker and the 8-arm PEGnor monomer, enabling both monomers to diffuse into the first layer. Therefore, it is reasonable to conclude that transport by diffusion is playing a dominant role in the observed increased interface thickness with soak time.

When a more tightly crosslinked hydrogel (i.e., Hydrogel III) was used to fabricate the first layer, the overall interface thickness followed a similar trend, but was smaller (e.g., initial interface thickness was \( \sim 70 \mu m \)). This observation is supported by a lower initial swelling response, shown in Fig. 1E where Hydrogel III swelled 2.5x its initial mass when placed in an aqueous solvent, and also had a lower diffusivity compared to Hydrogel II. Taken together, the thickness of the interface is dependent on both convective and diffusive transport.

Regardless of the transport mechanism, the transported monomers will subsequently be able to react with free norbornenes and/or thiols within the first layer and/or with each other to form a semi-interpenetrating network. Non-idealities that result during network formation will leave free reactive groups in the first layer enabling covalent bonds to form between the two layers (Dušek and Dušková-Smrčková, 2000; Elliott and Bowman, 2002; Elliott et al., 2004, 2003; Metters and Hubbell, 2005). Both the 8-arm PEGnor monomer and the PEGdt crosslinker were found to diffuse into the first layer, but the concentration of each monomer
via diffusive transport quickly and more so for the larger 8-arm PEGnor monomer. When the concentration of the 8-arm PEGnor drops below ~2 mM, a network is unable to form. Thus for diffusive transport, it can be reasoned that an interface forms due to covalent bonds and to a lesser extent by the formation of a semi-interpenetrating network. However, when convective transport dominates, the interface that forms is likely a combination of covalent bonds and semi-interpenetrating networks.

Interestingly, interface thickness did not vary with soak time when the first layer hydrogel was fabricated from Hydrogel IV, the stiffest hydrogel used in this study. However, the interface thickness at the initial time point was relatively large at ~140 μm. It was observed that Hydrogel IV, like Hydrogel II, undergoes significant swelling (~3.2x its initial mass) after polymerization to reach equilibrium, which can be seen in Fig. 1E. Thus, it is not surprising that the interface thickness at the initial time point was large for this hydrogel formulation, similar to the case with Hydrogel II. The lower diffusivity of the PEGdt and PEGnor, however, led to reduced diffusion into the hydrogel over the 20 min of soak time and as a result no further increase in the interface thickness was observed with longer soak times.

Taken together, these observations support the idea that for these PEG hydrogels, a rapid swelling response leads to the formation of an initial interface, which arises due to convective transport of the monomers into the first layer hydrogel. The initial thickness of the interface depends on how far the solvent concentration in the precursor solution is from the equilibrium-swollen volume of the hydrogel. The farther from the equilibrium-swollen volume, the larger the resulting initial interface will be. However, with longer soak times, the interface grows predominantly due to diffusion of the monomers into the previously polymerized hydrogel layer. By controlling the relative size of the monomers to that of the mesh size of the first layer, it is possible to further control the thickness of the interface and even prevent it from growing, at least at short soak times.

Mechanically, the bilayer hydrogels formed from sequential polymerization of thiol–ene photo-clickable monomers led to robust hydrogels that supported large compressive strains (up to 45%) with no observable signs of failure. The robust interface that links the two layers is attributed to a combination of covalent bonds and entanglements that form as the monomers are transported into the first layer and subsequently polymerized. The overall compressive modulus of the bilayer hydrogels was dependent on the relative modulus of the two layers, but did not appear to depend strongly on the interface thickness. When the layers were similar in crosslink density, the overall compressive modulus of the bilayer hydrogels was similar to the single layer hydrogels. When the layers had different moduli, the overall compressive modulus mirrored that of the soft layer at low strains, but at high strains exhibited a modulus that was in between the modulus of each layer. These observations follow composite theory of materials, which also predicts that a relatively thin interface (~10% or less of the total height) would not contribute appreciably to the overall compressive modulus (Chawla, 2012; Herakovich, 1997; Kinneberg et al., 2015).

The interfacial characteristics between adjacent layers within multi-layer hydrogels are critical to the successful distribution of strains, and prevention of interface failure, under compressive loading. From nano- and micro-mechanical measurements using both AFM and nanoindentation, respectively, we observed a gradient in mechanical properties across the interface of the Bilayer IV:I hydrogel. This observation indicates that after the first layer is polymerized, a gradient in crosslink density forms at its surface, which is consistent with the observed gradient in the fluorescence. We confirmed in a separate experiment that in single layer hydrogels a gradient appears at the free surface of the hydrogel due to swelling (DuPont et al., 2010). In addition, the temporal swelling behavior of hydrogels (Bell and Peppas, 1996; Brannon-Peppas and Peppas, 1991; Khare and Peppas, 1995; Martens and Anseth, 2000) will also contribute to the formation of a gradient. However, it is possible that an inverse gradient could form as monomers from the second layer are transported into the first layer and subsequently react, which will increase the local crosslink density. The latter appears to have less of an affect since a gradient from stiff to soft is observed. Overall, the combination led to a smooth, gradual transition in the modulus across the stiff-soft hydrogel interface.

The transition of mechanical properties across the interface of the Bilayer IV:I hydrogel was evaluated using a sharp probe via AFM, to evaluate changes in mechanical properties with nanometer spatial resolution. The mechanical interface between the stiff and soft layers in the bilayer hydrogel varied over a distance of approximately 70 μm, which is substantially smaller, by ~2-fold, than the interface thickness that was measured using confocal microscopy. This observation points to potential differences between the variation in chemistry and mechanical properties across the interface between two dissimilar hydrogels. While the PEG crosslinker molecules and to a lesser degree the 8-arm PEGnor monomers are transported farther into and react with the first layer, it is likely that not all contribute significantly to the mechanical properties. Thus, the interface thickness may be distinctly different with respect to chemistry and elastic modulus.

While AFM using a sharp probe is highly sensitive to mechanical property changes in a heterogeneous material, nanoindentation with a large sphere is better capable of accurately measuring material properties. To that end, we used nanoindentation with a large sphere to evaluate changes in the overall modulus of the bilayer hydrogel. When the layers had different moduli, the overall compressive modulus of the bilayer hydrogel mirrored that of the soft layer at low strains, but at high strains exhibited a modulus that was in between the modulus of each layer.
determining elastic moduli in hydrogels. Strains under sharp AFM probes are constant and related to the probe half-angle, yet are difficult to quantify exactly (Atkins and Tabor, 1965; Johnson, 1970). As hydrogels exhibit a nonlinear relationship between stress and strain (Anseth et al., 1996), moduli determined from AFM without knowing the corresponding strain is challenging. Comparably, probing with a large sphere enables facile determination of strains and maintains criteria needed for Hertzian contact mechanics (i.e., when small strains are maintained), and so permits evaluation of elastic moduli at known strain values. In this study, nanoindentation using a 100 μm cono-spherical probe produced relatively small, ~4% strain in the hydrogel volume compressed by the indenter probe. Moreover, commercial nanoindenter systems enable mechanical property (e.g., elastic modulus) mapping over large (centimeter) regions in comparison to limited range with AFM (max of ~200 μm in x- and y-direction). Evaluation of the Bilayer IV:1 hydrogel using nanoindentation revealed a gradual modulus gradient, over a distance of ~300 μm, between the two hydrogel layers with a transition region from >150 kPa, in Hydrogel IV, to ~60 kPa, in Hydrogel I. These moduli compare well with properties observed at comparable strains in unconfined compression testing of single layer hydrogels in this study. The wide distance indicating the interface of the bilayer hydrogel, as compared to the ~70 μm mechanical interface observed using AFM, results from the large contact radius (~32.4 μm) of the cono-spherical nanoindenter probe, where the material within roughly 3–5 times that of the contact radius contributes to the measured elastic modulus. Thus the mechanical interface revealed by nanoindentation is wider than that of AFM, where the latter possesses better spatial resolution by using a very small (~10 nm) sharp probe. Nonetheless, nanoindentation using a large probe is able to map the modulus gradient that exists between the two dissimilar hydrogels and provide more accurate modulus values. This gradient was also mirrored in the local strain measurements as observed by strain measurements in bead tracking experiments. The length-scale over which strain varied across the bilayer interface was even greater (~1000 μm in length) than the modulus gradient observed using nanoindentation. We attribute this finding to the experimental methods, where relatively large unconfined compressive strains (i.e., 15%) were applied to the entire hydrogel construct. As a result of these large strains, deformation of the hydrogel leads to lateral expansion (Kinneberg et al., 2015), and therefore extends the strains over a much larger distance. Importantly, the broad strain field observed indicates that the graded interface between the two dissimilar hydrogels enables effective strain transfer between stiff and soft layers. Moreover, the strain field is distributed far from the interface in the bilayer hydrogel and does not include a sharp increase at the interface of the two materials, which would indicate potential for interfacial failure.

There are several limitations that are important to note. Our findings point towards diffusion and convection as playing an important role in the formation of the interface of the bilayer hydrogels. It was, however, not possible to decouple these factors experimentally and therefore mathematical models may be necessary to differentiate them and identify their contribution to the formation of the interface. Another limitation of this study is the lack of nanomechanical evaluation of the other bilayer conditions. Due to the similarity in mechanics of the layers in Bilayer III:II and II:III the interface is difficult to probe, and the much larger interfaces formed at longer soak times would be challenging to capture by the narrow grids afforded by AFM. However, understanding the properties of such interfaces formed when the first layer is composed of larger mesh sizes like that of Hydrogel II or III would be paramount in further identifying the effects of swelling and diffusion on interface mechanics. Also as discussed, due to the non-linear relationship between stress and strain exhibited by hydrogels (Kinneberg et al., 2015), AFM likely overestimated the moduli of Hydrogels IV and I. Nanoindentation of soft materials like hydrogels requires large probes (i.e., due to limited ability of commercial nanoindenters to detect low force levels), and so lacks the spatial resolution of sharp AFM probes to directly evaluate modulus gradients. Yet nanoindentation can readily evaluate a functional (e.g., modulus) gradient across the bilayer interface in large millimeter-sized arrays and provides a realistic measure of elastic modulus of each hydrogel layer and the intermediate interfacial region. While each approach has limitations, we conclude that the combination of techniques used herein enables a comprehensive understanding of small length scale (<100 μm) interfacial regions at high spatial resolution (i.e., using AFM), moduli of single hydrogels and modulus gradients at interfaces (i.e., using nanoindentation) and behavior (via strain mapping) spanning large, millimeter-sized regions.

5. Conclusions

In this study, we characterized the interface that forms between two layers of a bilayer PEG hydrogel that is fabricated by a step-and-repeat process of photo-polymerized PEG monomers. We identified that the interface forms as a result of two phenomena: (a) convective transport of monomers, which occurs rapidly due to changes in swelling in the first layer when the second layer is applied and (b) diffusive transport of monomers, which occurs over longer time scales, and is dependent on the relative size of the monomers to the mesh size of the first layer. It is important to note that these phenomena are not specific to PEG hydrogels and therefore extend to any bilayer hydrogel that is formed by a step-and-repeat process. Mechanically, a robust interface forms as evident by no signs of failure (up to the strains tested) and the compressive modulus of the bilayer hydrogels mirrored that of the layers and not the interface. Nanomechanical analysis confirmed that a gradient in modulus forms across the interface and subsequently leads to a gradual transfer of strain across the entire bilayer hydrogel in unconfined compression. Overall, this work provides new insight into the mechanisms that control the formation of an interface and the resulting interfacial properties when multiple layers are sequentially polymerized. The ability to control the interface between sequentially polymerized hydrogels is important to interfacial tissue engineering applications where multi-layer hydrogels are employed.

Acknowledgments

Research reported in this publication was partially supported by the NSF Career Award DMR #0847390 (A.H.A., S.J.B.) and NSF Career Award CBET #1055989 (J.W., V.L.F.), NIH Pharmaceutical Biotechnology Training Fellowship TGM008732C (A.H.A.), Department of Education’s Graduate Assistantship in Areas of National Need P200A120063 (A.H.A.) and NIST Measurement Science & Engineering Fellowship 70NANB10H027 (A.H.A.). The authors acknowledge the NSF Major Research Instrumentation Award (NSF CBET#1338154) and the University of Colorado, Boulder CO, USA, for funding the combined Raman spectroscopy – nanoindenter system used in this work. Specific commercial equipment, instruments, and materials that are identified in this report are listed in order to adequately describe the experimental procedure and are not intended to imply endorsement or recommendation by the National Institute of Standards and Technology (NIST).