RESEARCH ARTICLE

Antiderivatives and integral representations of associated Legendre functions with $\mu = \pm \nu$

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We obtain antiderivatives and complex integral representations for associated Legendre functions and Ferrers functions (associated Legendre functions on-the-cut) of the first and second kind with degree and order equal to within a sign.

Keywords: Associated Legendre functions; Ferrers functions; Integral representations; Gauss hypergeometric function

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1. Introduction

Using analysis for fundamental solutions of Laplace’s equation on Riemannian manifolds of constant curvature, we have previously obtained antiderivatives and integral representations for associated Legendre and Ferrers functions of the second kind with degree and order equal to within a sign. For instance in Cohl (2011) \cite{2}, Theorem 1, we derived using the $d$-dimensional hypersphere with $d = 2, 3, 4, \ldots$, an antiderivative and an integral representation for the Ferrers function of the second kind with order equal to the negative degree. In Cohl \& Kalnins (2012) \cite{3}, Theorem 3.1, we derived using the $d$-dimensional hyperboloid model of hyperbolic geometry with $d = 2, 3, 4, \ldots$, an antiderivative and an integral representation for the associated Legendre function of the second kind with degree and order equal to each other. In \cite{2}, \cite{3}, the antiderivatives and integral representations were restricted to values of the degree and order $\nu$ such that $2\nu$ is an integer.

The purpose of this paper is to generalize the results presented in \cite{2}, \cite{3} for associated Legendre and Ferrers functions of the first and second kind, and to extend them such that the degree and order are no longer subject to the above restriction. Our extremely simple integral representations are consistent with known special values for associated Legendre and Ferrers functions of the first kind when the order is equal to the negative degree.

Throughout this paper we rely on the following definitions. Let $a_1, a_2, a_3, \ldots \in \mathbb{C}$, with $\mathbb{C}$ being the set of complex numbers. If $i, j \in \mathbb{Z}$ and $j < i$, then $\prod_{n=i}^{j} a_n = 1$. The set of natural numbers is given by $\mathbb{N} := \{1, 2, 3, \ldots\}$, the set $\mathbb{N}_0 := \{0, 1, 2, \ldots\} = \mathbb{N} \cup \{0\}$, and $\mathbb{Z} := \{0, \pm 1, \pm 2, \ldots\}$.

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2. Associated Legendre functions of the first and second kind

As is the common convention for associated Legendre functions \[1\] (8.1.1), for any expression of the form \((z^2 - 1)^\alpha\), read this as

\[
(z^2 - 1)^\alpha := (z + 1)^\alpha (z - 1)^\alpha,
\]

for any fixed \(\alpha \in \mathbb{C}\) and \(z \in \mathbb{C} \setminus (-\infty, 1]\).

**Theorem 2.1** Let \(z \in \mathbb{C} \setminus (-\infty, 1]\), \(\nu \in \mathbb{C} \setminus \{-\frac{1}{2}, -\frac{3}{2}, -\frac{5}{2}, \ldots\}\), and \(C\) be a constant.

Then we have the following antiderivative

\[
\int \frac{dz}{(z^2 - 1)^{\nu + 1}} = \frac{-1}{(2\nu + 1)z^{2\nu + 1}} \mathbf{2F}_1 \left( \frac{\nu + 1}{2}, \frac{\nu + 1}{2} ; \frac{1}{z^2} \right) + C
\]

\[
= \frac{-2^{-\nu} e^{-i\pi \nu}}{\Gamma(\nu + 1)(z^2 - 1)^{\nu/2}} Q_\nu(z) + C. \tag{1}
\]

In the above expression, the Gauss hypergeometric function \(\mathbf{2F}_1\) can be defined in terms of the following infinite series

\[
\mathbf{2F}_1 \left( a, b ; c ; z \right) := \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n} \frac{z^n}{n!},
\]

(see (2.1.5) in Andrews, Askey & Roy 1999) and elsewhere on \(z \in \mathbb{C} \setminus (1, \infty)\) by analytic continuation. The Pochhammer symbol (rising factorial) \((\cdot)_n : \mathbb{C} \to \mathbb{C}\) is defined by

\[
(z)_n := \prod_{i=1}^{n} (z + i - 1),
\]

where \(n \in \mathbb{N}_0\). The associated Legendre function of the second kind \(Q_\nu(z) : \mathbb{C} \setminus (-\infty, 1] \to \mathbb{C}\) can be defined in terms of the Gauss hypergeometric function for \(\nu + \mu \notin -\mathbb{N}\), as (Olver et al. (2010) \[2\], (14.3.7) and section 14.21)

\[
Q_\nu(z) := \frac{\sqrt{\pi} e^{i\pi \mu} \Gamma(\nu + \mu + 1)(z^2 - 1)^{\mu/2}}{2^{\nu+1} \Gamma(\nu + \frac{3}{2}) z^{\nu+\mu+1}} \mathbf{2F}_1 \left( \frac{\nu+\mu+1}{2}, \frac{\nu+\mu+1}{2} ; \frac{1}{z^2} \right), \tag{2}
\]

for \(|z| > 1\) and by analytic continuation of the Gauss hypergeometric function elsewhere on \(z \in \mathbb{C} \setminus (-\infty, 1]\).

**Proof.** The antiderivative (1) is verified as follows. By using

\[
\frac{d}{dz} \mathbf{2F}_1 \left( a, b ; c ; z \right) = \frac{ab}{c} \mathbf{2F}_1 \left( a + 1, b + 1 ; c + 1 ; z \right) \tag{3}
\]
(see (15.5.1) in Olver et al. (2010) [7]), and the chain rule, it follows that
\[
\frac{1}{(2\nu + 1)} \frac{d}{dz} \frac{1}{z^{2\nu+1}} \, _2F_1\left(\frac{\nu + \frac{1}{2}, \nu + 1}{\nu + \frac{3}{2}} \right) = \frac{-1}{z^{2\nu+2}} \, _2F_1\left(\frac{\nu + \frac{1}{2}, \nu + 1}{\nu + \frac{3}{2}} \right)
\]
\[
- \frac{2(\nu + 1)}{(2\nu + 3)z^{2\nu+1}} \, _2F_1\left(\frac{\nu + \frac{3}{2}, \nu + 2}{\nu + \frac{5}{2}} \right).
\]
(4)

The second hypergeometric function on the right-hand side can be simplified using Gauss’ relations for contiguous hypergeometric functions, namely
\[
z \, _2F_1\left(\frac{a + 1, b + 1}{c + 1} \right) = \frac{c}{a - b} \left[ _2F_1\left(\frac{a, b + 1}{c} \right) - _2F_1\left(\frac{a + 1, b}{c} \right) \right]
\]
(5)
(see p. 58 in Erdélyi et al. (1981) [4]), and
\[
_2F_1\left(\frac{a, b + 1}{c} \right) = \frac{b - a}{b} \, _2F_1\left(\frac{a, b}{c} \right) + \frac{a}{b} \, _2F_1\left(\frac{a + 1, b}{c} \right)
\]
(6)
(see (15.5.12) in Olver et al. (2010) [7]). After simplification and utilization of
\[
_2F_1\left(\frac{a, b}{c} \right) = (1 - z)^{-a}
\]
(see (15.4.6) in Olver et al. (2010) [7]), the right-hand side of (11) simplifies to $1/(z^2 - 1)^{\nu+1}$. The equality of the Gauss hypergeometric function in the antiderivative (11), in terms of the associated Legendre function $Q_\nu$, follows directly from (12).

A straightforward consequence of the antiderivative (11) is the following integral representation for the associated Legendre function with degree and order equal, namely
\[
Q_\nu(z) = 2^\nu \Gamma(\nu + 1) e^{i\nu \pi} (z^2 - 1)^{\nu/2} \int_{z}^{\infty} \frac{dw}{(w^2 - 1)^{\nu+1}},
\]
(7)
where Re $\nu > -\frac{1}{2}$. Using the negative order relation for associated Legendre functions of the second kind [7, (14.9.14)]
\[
Q_{\nu}^{-\mu}(z) = e^{-2i\pi\mu} \frac{\Gamma(\nu - \mu + 1)}{\Gamma(\nu + \mu + 1)} Q_\nu^\mu(z),
\]
we derive
\[
Q_{\nu}^{-\nu}(z) = \frac{\sqrt{\pi} e^{-i\nu \pi} (z^2 - 1)^{\nu/2}}{2^\nu \Gamma(\nu + 1/2)} \int_{z}^{\infty} \frac{dw}{(w^2 - 1)^{\nu+1}}.
\]
Using the Whipple relation for associated Legendre functions (cf. [7, (14.9.16)])
\[
Q_\nu(z) = \sqrt{\frac{\pi}{2}} \Gamma(2\nu + 1)(z^2 - 1)^{-1/4} e^{i\nu \pi} P_{-\nu-1/2}^{-\nu-1/2} \left( \frac{z}{\sqrt{z^2 - 1}} \right),
\]
where $P_{\mu}(x) = \frac{d}{dx} P_{\mu+1}(x)$ is the Legendre polynomial. (13)
where the associated Legendre function of the first kind $P_{\nu}^{\mu} : C \setminus (-\infty, 1] \rightarrow C$ is defined in [7, (14.3.6)], we obtain

$$P_{\nu}^{\mu}(z) = -\frac{2^{\nu+1}}{\pi} \Gamma(\nu + 1) \sin(\nu \pi) (z^2 - 1)^{\nu/2} \int_{z/\sqrt{z^2 - 1}}^{\infty} \frac{dw}{(w^2 - 1)^{-\nu+1/2}}.$$  

The above expression is equivalent to

$$P_{\nu}^{\mu}(z) = \frac{2^\nu \Gamma(\nu + 1/2)}{\sqrt{\pi}} (z^2 - 1)^{\nu/2}$$

$$+ \frac{2^{\nu+1} \Gamma(\nu + 1)(z^2 - 1)^{\nu/2}}{\pi} \sin(\nu \pi) \int_{z}^{\infty} \frac{dw}{(w^2 - 1)^{\nu+1}}, \quad (8)$$

where $\text{Re} \, \nu > -\frac{1}{2}$. An interesting definite integral follows from the behavior of the above integral representation near the singularity at $z = 1$, namely

$$\int_{1}^{\infty} (w^2 - 1)^{\nu-1} dw = \frac{\Gamma(\nu) \Gamma\left(\frac{1}{2} - \nu\right)}{2\sqrt{\pi}},$$

for $0 < \text{Re} \, \nu < \frac{1}{2}$. Using the negative order relation for associated Legendre functions of the first kind (cf. [7, (14.9.12)]), one has

$$P_{-\nu}^{\nu}(z) = \frac{1}{\Gamma(2\nu + 1)} \left[ P_{\nu}^{\nu}(z) - \frac{2}{\pi} e^{-i\nu \pi} \sin(\nu \pi) Q_{\nu}^{\nu}(z) \right]. \quad (9)$$

After replacement of (7) and (8) in (9) we obtain

$$P_{-\nu}^{\nu}(z) = \frac{\sin(\nu \pi) (z^2 - 1)^{\nu/2}}{\sqrt{\pi} 2^{\nu-1} \Gamma(\nu + 1/2)} \left[ \int_{z}^{\infty} \frac{dw}{(w^2 - 1)^{-\nu+1/2}} + \int_{z}^{\infty} \frac{dw}{(w^2 - 1)^{\nu+1}} \right],$$

which reduces to the special value [7, (14.5.19)]

$$P_{-\nu}^{\nu}(z) = \frac{(z^2 - 1)^{\nu/2}}{2^\nu \Gamma(\nu + 1)}.$$

3. Ferrers functions of the first and second kind

**Theorem 3.1** Let $x \in (-1, 1)$, $\nu \in C$, and $C$ be a constant. Then we have the following antiderivative

$$\int \frac{dx}{(1 - x^2)^{\nu+1}} = x_2 F_1 \left( \frac{1}{2}, \nu + 1; \frac{3}{2}; x^2 \right) + C = \frac{2^\nu \Gamma(\nu + 1/2)}{\sqrt{\pi} (1 - x^2)^{\nu/2}} Q_{\nu}^{\nu}(x) + C.$$
and for $\mu = -\nu$ we have
\[
Q^{-\nu}_{\nu}(x) = \frac{\sqrt{\pi} x (1 - x^2)^{\nu/2}}{2^\nu \Gamma(\nu + \frac{1}{2})} {}_2F_1\left(\frac{1}{2}, \nu + 1; \nu; x^2\right).
\] (10)

**Proof.** The Gauss hypergeometric function in the antiderivative follows using (3), (5), (6), as in the proof of Theorem 2.1, with the Ferrers function following directly using (10).

The following integral representation for the Ferrers function of the second kind is an obvious consequence of Theorem 3.1, namely
\[
Q^{-\nu}_{\nu}(x) = \sqrt{\pi} \frac{(1 - x^2)^{\nu/2}}{2^\nu \Gamma(\nu + \frac{1}{2})} \int_0^x \frac{dw}{(1 - w^2)^{\nu+1}}.
\] (11)

A definite-integral result near the singularity at $x = 1$ follows using (10), (11), and Gauss’s sum \[7, (15.4.20)\], namely
\[
P^{-\nu}_{\nu}(x) = \frac{1 - x^2)^{\nu/2}}{2^\nu \Gamma(\nu + 1)}
\]
for $\text{Re}(c - a - b) > 0$, namely
\[
\int_0^1 (1 - w^2)^{\nu-1} dw = \frac{\sqrt{\pi} \Gamma(\nu)}{2^\nu \Gamma(\nu + \frac{1}{2})},
\]
for $\text{Re}\nu > 0$. The well-known special value (see \[7, (14.5.18)\])
\[
P^{-\nu}_{\nu}(x) \Gamma(c-a-b) = \frac{1}{\Gamma(c-a)\Gamma(c-b)}
\]
where $P^\mu_{\nu} : (-1, 1) \to \mathbb{C}$ is the Ferrers function of the first kind (associated Legendre function of the first kind on-the-cut) defined in \[7, (14.3.1)\], in conjunction with \[7, (8.737.1)\], yields the following integral representation
\[
P^\mu_{\nu}(x) = \frac{2^\nu (1 - x^2)^{\nu/2}}{\sqrt{\pi}} \left[ \Gamma(\nu + \frac{1}{2}) \cos(\nu \pi) + \frac{2\Gamma(\nu + 1)}{\sqrt{\pi}} \sin(\nu \pi) \int_0^x \frac{dw}{(1 - w^2)^{\nu+1}} \right].
\]

Finally using the negative order relation for the Ferrers functions of the second kind \[6, p. 170\], namely
\[
Q^\nu_{-\mu}(x) = \frac{\Gamma(\nu - \mu + 1)}{\Gamma(\nu + \mu + 1)} \left[ \cos(\mu \pi)Q^\nu_{\nu}(x) + \frac{\pi}{2} \sin(\mu \pi)P^\mu_{\nu}(x) \right],
\]
we have the following integral representation of the Ferrers function of the second kind with degree and order equal
\[
Q^\nu_{\nu}(x) = -2^{\nu-1} \sqrt{\pi} \Gamma(\nu + \frac{1}{2}) \sin(\nu \pi)(1 - x^2)^{\nu/2}
\]
\[
+ 2^\nu \Gamma(\nu + 1) \cos(\nu \pi)(1 - x^2)^{\nu/2} \int_0^x \frac{dw}{(1 - w^2)^{\nu+1}}.
\]
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References


